

Pockels & Kerr Effects.

$\vec{E} = (E_1, E_2, E_3)$ is applied to a crystal
 $\Rightarrow \vec{\eta}$ permeability tensor is altered.

$\vec{\eta} = \epsilon_0 \vec{\epsilon}^{-1}$ $\eta_{ij} = \eta_{ji}$

Apply $\vec{E} \Rightarrow \eta_{ij} = \eta_{ij}(\vec{E})$.

$\eta_{ij}(\vec{E}) = \eta_{ij} + \sum_k r_{ijk} E_k + \sum_{kl} s_{ijkl} E_k E_l, \quad i, j, k, l = 1, 2, 3, \dots$

$\eta_{ij} = \eta_{ij}(0) ; r_{ijk} = \frac{\partial \eta_{ij}}{\partial E_k} \Big|_{\vec{E}=0} \quad s_{ijkl} = \frac{1}{2} \frac{\partial^2 \eta_{ij}}{\partial E_k \partial E_l} \Big|_{\vec{E}=0}$

\rightarrow generalized form of $\eta(\vec{E}) = \eta + rE + sE^2$

r replaced by $3^3 = 27$ coefficients $\{r_{ijk}\}$.

s replaced by $3^4 = 81$ coefficients $\{s_{ijkl}\}$.

r_{ijk} - Pockels coefficients - linear electro-optic. 3rd rank tensor

s_{ijkl} - Kerr coefficients - quadratic electro-optic coeff. 4th rank tensor

Symmetry. $\vec{\eta}$ - symmetric $r_{ijk} \quad s_{ijkl}$ invariant under permutations of i and j .

$r_{ijk} = r_{jik} \quad s_{ijkl} = s_{jike}$

in addition: $s_{ijkl} = s_{ijlk}$

ij permutations \Rightarrow 6 elements eg: 11, 12, 13, 22, 23, 33

$kl \Rightarrow$ 6 elements. $\rightarrow r_{ijk}$ has 6×3 elements

$s_{ijkl} = 6 \times 6$ elements.

We rename $(i, j), i, j = 1, 2, 3 \rightarrow$ single index
 $I = 1, 2, 3, \dots, 6$

$r_{ijk} \rightarrow r_{IK} \quad S_{ijkl} \rightarrow S_{IK} \quad K = 1, 2, 3, \dots, 6$

example $r_{12K} \rightarrow r_{6K} \quad S_{1231} \rightarrow S_{65}$

From Table:

j	i=1	2	3
1	1	6	5
2	6	2	4
3	5	4	3

$(i, j) = (3, 2) \rightarrow I = 4$

r (3rd rank tensor) $\rightarrow 6 \times 3$ matrix
 S (4th rank tensor) $\rightarrow 6 \times 6$ matrix.

Crystal symmetry.

Crystal symmetry imposes more constraints on r & S
 \rightarrow some entries must be equal, some zero, opposite in sign or some other relat.

Example Pockels Effect.

$$\begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ r_{41} & 0 & 0 \\ 0 & r_{41} & 0 \\ 0 & 0 & r_{41} \end{bmatrix}$$

Cubic $\bar{4}3m$
 GaAs, CdTe, InAs.

$$\begin{bmatrix} r_{41} \\ r_{41} \\ r_{63} \end{bmatrix}$$

tetragonal $\bar{4}2m$
 (KDP, ADP)

$$\begin{bmatrix} 0 & -r_{22} & r_{13} \\ 0 & r_{22} & r_{13} \\ 0 & 0 & r_{33} \\ 0 & r_{51} & 0 \\ r_{51} & 0 & 0 \\ -r_{22} & 0 & 0 \end{bmatrix}$$

Trigonal $3m$
 (LiNbO₃, LiTaO₃).

Kerr Coefficients. S_{ijk} for an isotropic medium.

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$$\begin{bmatrix} S_{11} & S_{12} & S_{12} & 0 & 0 & 0 \\ S_{21} & S_{11} & S_{12} & 0 & 0 & 0 \\ S_{12} & S_{12} & S_{11} & 0 & 0 & 0 \\ 0 & 0 & 0 & S_{44} & 0 & 0 \\ 0 & 0 & 0 & 0 & S_{44} & 0 \\ 0 & 0 & 0 & 0 & 0 & S_{44} \end{bmatrix}$$

$$S_{44} = \frac{S_{11} - S_{12}}{2}$$

Pockels Effect. $\vec{E} = (E_1, E_2, E_3)$

Determine optical properties.

- ① Find principal axes and principal refractive indices n_1, n_2 and n_3 in absence of \vec{E} .
- ② Find $\{r_{ijk}\}$ by using r_{ik} along with contraction rule relating i, j to \mathbb{I} .
- ③ Determine impermeability tensor using.
$$\eta_{ij}(\vec{E}) = \eta_{ij}(0) + \sum_k r_{ijk} E_k$$

where $\eta_{ij}(0)$ is a diagonal matrix with elements $\frac{1}{n_1^2}, \frac{1}{n_2^2}, \frac{1}{n_3^2}$

- ④ Write equation for modified index ellipsoid:

$$\sum_{ij} \eta_{ij}(\vec{E}) x_i x_j = 1$$

- ⑤ Determine principal axes of modified index ellipsoid by diagonalizing the matrix $\eta_{ij}(\vec{E})$ and find the corresponding principal refractive indices $n_1(\vec{E}), n_2(\vec{E})$ & $n_3(\vec{E})$.

- ⑥ Given direction of light propagation, \vec{k} , find the normal modes and their associated refractive indices using the index ellipsoid.

Index ellipsoid

$$\vec{\eta} = \begin{bmatrix} 1/n_1^2 & 0 & 0 \\ 0 & 1/n_2^2 & 0 \\ 0 & 0 & 1/n_3^2 \end{bmatrix}$$

represent graphically as.

$$\sum \eta_{ij} x_i x_j = 1.$$

Write as equation: $\eta_{11} x_1^2 + \eta_{12} x_1 x_2 + \eta_{13} x_1 x_3 + \eta_{21} x_1 x_2 + \eta_{22} x_2^2 + \eta_{23} x_2 x_3 + \eta_{31} x_3 x_1 + \eta_{32} x_3 x_2 + \eta_{33} x_3^2 = 1$

but $\eta_{11} = \frac{1}{n_1^2}$, $\eta_{22} = \frac{1}{n_2^2}$, $\eta_{33} = \frac{1}{n_3^2}$ all others zero.

$$\Rightarrow \frac{x_1^2}{n_1^2} + \frac{x_2^2}{n_2^2} + \frac{x_3^2}{n_3^2} = 1.$$

or
$$\boxed{\frac{x^2}{n_1^2} + \frac{y^2}{n_2^2} + \frac{z^2}{n_3^2} = 1.}$$

equation for ellipsoid!

n_1, n_2, n_3 are refractive indices.

Pockels Effect in anisotropic media:

$$\eta_{ij} = \eta_{ij}(0) + \sum_k \gamma_{ijk} E_k.$$

$$\begin{bmatrix} \eta_{11} & \eta_{12} & \eta_{13} \\ \eta_{21} & \eta_{22} & \eta_{23} \\ \eta_{31} & \eta_{32} & \eta_{33} \end{bmatrix} = \begin{bmatrix} 1/n_1^2 & 0 & 0 \\ 0 & 1/n_2^2 & 0 \\ 0 & 0 & 1/n_3^2 \end{bmatrix} + \sum_k \gamma_{ijk} E_k.$$

Review of concepts - clarifications.



TEM plane waves: $\vec{E}(z, t) = \text{Re} \left\{ \vec{A} \exp \left[jz\pi\nu \left(t - \frac{z}{c} \right) \right] \right\}$

we know always there.

and $\vec{A} = A_x \hat{x} + A_y \hat{y}$ A_x & A_y are complex.

Jones vector: $\mathcal{J} = \begin{bmatrix} A_x \\ A_y \end{bmatrix}$ completely describes wave!

$\vec{D} = \epsilon \vec{E} + \vec{P}$ $D_i = \sum_j \epsilon_{ij} E_j$ anisotropic medium.

or $\vec{D} = \overleftrightarrow{\epsilon} \vec{E}$.

example.

$$D_x = \epsilon_{xx} E_x + \epsilon_{xy} E_y + \epsilon_{xz} E_z$$

$$D_y = \epsilon_{yx} E_x + \epsilon_{yy} E_y + \epsilon_{yz} E_z$$

$$D_z = \epsilon_{zx} E_x + \epsilon_{zy} E_y + \epsilon_{zz} E_z$$

or $\begin{bmatrix} D_x \\ D_y \\ D_z \end{bmatrix} = \begin{bmatrix} \epsilon_{xx} & \epsilon_{xy} & \epsilon_{xz} \\ \epsilon_{yx} & \epsilon_{yy} & \epsilon_{yz} \\ \epsilon_{zx} & \epsilon_{zy} & \epsilon_{zz} \end{bmatrix} \begin{bmatrix} E_x \\ E_y \\ E_z \end{bmatrix}$

permittivity tensor.

Rewrite $\overleftrightarrow{\eta} = \epsilon_0 \overleftrightarrow{\epsilon}^{-1}$: impermeability tensor.

$\vec{E} = \frac{\overleftrightarrow{\eta}}{\epsilon_0} \vec{D}$ o.k. ? or $\vec{E} = \frac{1}{\epsilon_0} \sum \eta_{ij} D_j$

In addition: we define $n_1 = \left(\frac{\epsilon_1}{\epsilon_0} \right)^{1/2}$, $n_2 = \left(\frac{\epsilon_2}{\epsilon_0} \right)^{1/2}$, $n_3 = \left(\frac{\epsilon_3}{\epsilon_0} \right)^{1/2}$

$$\sum_k \gamma_{ijk} E_k = ?$$

$i, j, k = 1, 2, 3$. 27 coefficients.

How many total elements. i, j are indices \Rightarrow 9 final values.

$$b_{21} \left[\begin{array}{l} \gamma_{111} E_1 + \gamma_{112} E_2 + \gamma_{113} E_3, \\ \gamma_{211} E_1 + \gamma_{212} E_2 + \gamma_{213} E_3, \\ \gamma_{311} E_1 + \gamma_{312} E_2 + \gamma_{313} E_3 \end{array} \right]$$

b_{12}

But from symmetry: $\gamma_{ijk} = \gamma_{jki}$.

$$\Rightarrow \gamma_{121} = \gamma_{211}$$

$$\therefore b_2 = \gamma_{121} E_1 + \gamma_{122} E_2 + \gamma_{123} E_3 = \gamma_{211} E_1 + \gamma_{212} E_2 + \gamma_{213} E_3 = b_{21}$$

Trigonal 3m Crystals. $n_1 = n_2 = n_o$, $n_3 = n_e$ r given earlier. 42/

$\vec{E} = (0, 0, E)$ Electric field along optic axis.

$$\eta_{ij}(E_z) = \eta_{ij}(0) + \gamma_{ijz} E_z.$$

$$\eta_I(E_z) = \eta_I(0) + \gamma_{Iz} E_z. \quad z=3.$$

only $I=1,2,3$ are non zero γ_{13}, γ_{33} .

$$\eta_1 = \eta_1(0) + \gamma_{13} E_3. \quad \text{and } \gamma_{23} = \gamma_{13}.$$

$$\eta_3 = \eta_3(0) + \gamma_{33} E_3.$$

$$\eta_2 = \eta_2(0) + \gamma_{23} E_3. \quad (\gamma_{23} = \gamma_{13})$$

$$\eta_1 = \frac{1}{n_1^2} \quad \eta_3 = \frac{1}{n_3^2}$$

$$\frac{1}{n_3^2(E)} = \frac{1}{n_3^2} + \gamma_{33} E_3 \quad \frac{1}{n_1^2(E)} = \frac{1}{n_1^2} + \gamma_{13} E_3$$

$$\left(\frac{1}{n_o^2} + \gamma_{13} E \right) (x_1^2 + x_2^2) + \left(\frac{1}{n_e^2} + \gamma_{33} E \right) x_3^2 = 1.$$

$$\frac{1}{n_o^2(E)} = \frac{1}{n_o^2} + \gamma_{13} E$$

$$\frac{1}{n_e^2(E)} = \frac{1}{n_e^2} + \gamma_{33} E$$

$\gamma_{13} E$ & $\gamma_{33} E$ are generally small $\Rightarrow (1+\Delta)^{-1/2} \approx 1 - \frac{1}{2}\Delta$. $|\Delta|$ small.

$$\Rightarrow n_o(E) \approx n_o - \frac{1}{2} n_o^3 \gamma_{13} E$$

$$n_e(E) \approx n_e - \frac{1}{2} n_e^3 \gamma_{33} E$$

} crystal remains uniaxial with same principal axes.

Tetragonal 42m Crystals (KDP & ADP).

$\gamma_{41}, \gamma_{52} = \gamma_{41}, \gamma_{63}$ only non zero elements.

$$\vec{E} = (0, 0, E).$$

$$\eta_4(E) = \eta_4(0) + \gamma_{41} E_1 = \eta_4(0) = \eta_{32}(0) = \eta_{23}(0) = 0$$

$$\eta_5(E) = \eta_5(0) + \gamma_{51} E_1 = \eta_5(0) = \eta_{31}(0) = \eta_{23}(0) = 0$$

$$\eta_6(E) = \eta_6(0) + \gamma_{63} E_3 = \eta_6(0) + \gamma_{63} E = \gamma_{63} E$$

$$\eta_{42}(E) = \eta_{42}(0) + \gamma_{63} E = \gamma_{63} E$$

$$n_2(E) = n_{20}(0) + \gamma_{63} E = \gamma_{63} E$$

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$$n_1(E) = \frac{1}{n_1^2} = \frac{1}{n_0^2}$$

$$n_2(E) = \frac{1}{n_2^2} = \frac{1}{n_0^2}$$

$$n_3(E) = \frac{1}{n_3^2}$$

$$\Rightarrow \frac{x_1^2 + x_2^2}{n_0^2} + \frac{x_3^2}{n_0^2} + 2\gamma_{63} E x_1 x_2 = 1.$$

Can be rewritten as (by diagonalizing \vec{n} matrix).

$$\frac{u_1^2}{n_1^2(E)} + \frac{u_2^2}{n_2^2(E)} + \frac{u_3^2}{n_3^2(E)} = 1.$$

$$\frac{1}{n_1^2(E)} = \frac{1}{n_0^2} + \gamma_{63} E$$

$$\frac{1}{n_2^2(E)} = \frac{1}{n_0^2} - \gamma_{63} E$$

$$n_3(E) = n_e.$$

uniaxial crystal
 \Rightarrow biaxial crystal.

$$(1+\Delta)^{-1/2} \approx 1 - \frac{1}{2}\Delta$$

$$\Rightarrow n_1(E) = n_0 - \frac{1}{2} n_0^3 \gamma_{63} E$$

$$n_2(E) = n_0 + \frac{1}{2} n_0^3 \gamma_{63} E$$

$$n_3(E) = n_e.$$

Cubic Crystals $\bar{4}3m$. (GaAs, CdTe, InAs). $n_1 = n_2 = n_3 = n$
 isotropic. $\gamma_{41} = \gamma_{52} = \gamma_{63}$.

$$\vec{E} = (0, 0, E).$$

$$n_1 = n_1(0)$$

$$n_2 = n_2(0)$$

$$n_3 = n_3(0)$$

$$n_4 = \gamma_{41} E_1 = 0$$

$$n_5 = \gamma_{52} E_2 = 0$$

$$n_6 = \gamma_{63} E_3 = \gamma_{41} E$$

$$\Rightarrow n_{21} = n_{12} = \gamma_{41} E$$

$$\frac{x_1^2 + x_2^2 + x_3^2}{n^2} + 2r_{41}E x_1 x_2 = 1$$

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● new axes rotated by 45° about Z axis

$$\left. \begin{aligned} n_1(E) &\approx n - \frac{1}{2}n^3 r_{41} E \\ n_2(E) &\approx n + \frac{1}{2}n^3 r_{41} E \\ n_3(E) &\approx n \end{aligned} \right\} \begin{array}{l} \text{biaxial crystal.} \\ \Rightarrow \text{uniaxial} \\ \rightarrow \text{biaxial.} \end{array}$$

Notice in all of the above, the crystals principal axes do not change. as the electric field E increases
 \Rightarrow Normal modes stay the same, refractive indices change
 \Rightarrow Crystals can be used as a phase modulator, wave retarder or intensity modulator.

Kerr Effect $n_{ij}(E) = n_{ij}(0) + \sum_k r_{ijkl} E_k E_l$

Isotropic medium: $\vec{E} = (0, 0, E)$

$$\left(\frac{1}{n^2} + S_{12} E^2 \right) (x_1^2 + x_2^2) + \left(\frac{1}{n^2} + S_{11} E^2 \right) x_3^2 = 1$$

$$\frac{1}{n_o^2(E)} = \frac{1}{n^2} + S_{12} E^2$$

$$\frac{1}{n_e^2(E)} = \frac{1}{n^2} + S_{11} E^2$$

or $\left. \begin{aligned} n_o(E) &\approx n - \frac{1}{2}n^3 S_{12} E^2 \\ n_e(E) &\approx n - \frac{1}{2}n^3 S_{11} E^2 \end{aligned} \right\} \text{quadratic decreasing functions of } E.$

Modulators. from Pockels cells - can also be Kerr cells.

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Phase modulators. $n(E) \approx n - \frac{1}{2} r n^3 E$ n and r
refractive index \uparrow Pockels coefficient

$$E = \frac{V}{d} \quad \phi = \phi_0 - \pi \frac{V}{V_\pi} \quad \text{where } \phi_0 = \frac{2\pi n L}{\lambda_0}$$

$$V_\pi = \frac{d \lambda_0}{L r n^3} \quad \text{: half wave voltage. } n \text{ and } r$$

can be determined as follows.

Trigonal 3m Crystal. $\vec{E} = (0, 0, E)$. \rightarrow stays uniaxial.

$$n_o(E) \approx n_o - \frac{1}{2} n_o^3 r_{13} E$$

$$n_e(E) \approx n_e - \frac{1}{2} n_e^2 r_{33} E$$

Longitudinal modulator: $\vec{k} = k \hat{z} = (0, 0, k)$ parallel to E
 $\Rightarrow n = n_o \quad r = r_{13}, \quad d = L$

$\text{LiNbO}_3 \quad r_{13} = 9.6 \text{ pm/V}, \quad n_o = 2.3, \quad \lambda_0 = 633 \text{ nm (HeNe Laser)}$

$V_\pi = 5.41 \text{ kV}, \quad 5.41 \text{ kV is required for } \pi \text{ phase change.}$

Transverse modulator. $\vec{k} = (k, 0, 0)$ x-direction.

and is polarized in z-direction

$$\Rightarrow n = n_e \quad r = r_{33}. \quad d \neq L.$$

$\text{LiNbO}_3 \quad r_{33} = 30.9 \text{ pm/V} \quad n_e = 2.2, \quad V_\pi = 1.9 \left(\frac{d}{L}\right) \cdot \text{kV}$

if $\frac{d}{L} = 0.1 \Rightarrow V_\pi \approx 190 \text{ V}$ - much smaller than longitudinal case.

Intensity Modulator.

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$$\Gamma = \Gamma_0 - \pi \frac{V}{V_{\pi}}$$

$$\Gamma_0 = \frac{2\pi (n_1 - n_2) L}{\lambda_0}$$

$$V_{\pi} = \frac{(d/L) \lambda_0}{r_1 n_1^3 - r_2 n_2^3}$$

Cell between crossed polarizers \Rightarrow intensity modulator.

Example. Tetragonal $\bar{4}2m$ Crystal (KDP & ADP).

$$\left. \begin{array}{l} \vec{E} = (0, 0, E) \\ \vec{k} = (0, 0, k) \end{array} \right\} \Rightarrow \text{Normal modes}$$

$$n_1 = n_2 = n_0$$

$$r_1 = r_{63}$$

$$r_2 = -r_{63}$$

recall

$$\left\{ \begin{array}{l} n_1(E) = n_0 - \frac{1}{2} n_0^3 r_{63} E \\ n_2(E) = n_0 + \frac{1}{2} n_0^3 r_{63} E \\ n_3(E) = n_e \end{array} \right.$$

If $d/L \Rightarrow 1$ with $n_1 = n_2 = n_0 \Rightarrow \Gamma_0 = 0$.

$$V_{\pi} = \frac{\lambda_0}{2 r_{63} n_0^3}$$

KDP at $\lambda_0 = 633 \text{ nm}$

$$V_{\pi} = 8.4 \text{ kV.}$$